

Impact of Space-Dependent Thermal Conductivity on Heat Transfer in a Vertical Annulus with Asymmetric Surface Heating

Michael O Oni^{1*}, Taiwo S Yusuf¹, Olaife H Adebayo¹ and Luqman A Azeez²

¹Department of Mathematics, Ahmadu Bello University, Nigeria

²Federal College of Education, Nigeria

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***Corresponding author:** Michael O Oni, Department of Mathematics, Ahmadu Bello University, Zaria, Nigeria

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Abstract

This technical brief investigates the impact of space dependent thermal conductivity on temperature distributions and rate of heat transfer of some fluids in a porous annulus. The heat transport equation is derived with a linear thermal conductivity model. Exact solution of the energy equation is obtained after using suitable dimensionless parameters. Results indicate that increase in coefficient of thermal conductivity leads to a corresponding increase in fluid temperature as well as Nusselt number. In addition, fluids with higher Prandtl number are found to be a good insulators.

Keywords: Space-dependent; Thermal conductivity; Vertical annulus; Asymmetric heating; Prandtl number; Nusselt number

Introduction

Over the past two centuries, thermal conductivity (often denoted k , λ , or κ) is the property of a material to conduct heat. It is evaluated primarily in terms of Fourier's law for heat conduction. Materials of high thermal conductivity transfer heat faster than materials with lower thermal conductivity. Correspondingly, materials of high thermal conductivity are widely used in heat sink applications and materials of low thermal conductivity are used as thermal insulation. The thermal conductivity of a material may depend on temperature. Thermal conductivity is important in material science research, electronics, building insulation and related fields, especially where high operating temperatures are achieved. The impact of temperature dependent thermal conductivity has been studied over decades and interesting results have been obtained. Prasad et al. [1] studied the effects of variable fluid properties on the hydromagnetic flow and heat transfer over a non-linearly stretching sheet. Later, Animasaun [2] investigated the effect of thermophoresis, variable viscosity and thermal conductivity on free convective heat and mass transfer of non-darcian MHD dissipative casson fluid flow with suction and order of chemical reaction. He concluded that viscosity parameter as well as thermal conductivity parameter decreases fluid temperature. Other works on temperature dependent thermal conductivity can be found in [3-6].

Understanding heat transfer with space dependent thermal conductivity induced by temperature difference at the surfaces of the cylinders would give better outcome in laboratories experiment, material science, polymer factories and florescent bulbs industries. Liu and Yang [7] studied length-dependent thermal conductivity of single extended polymer chains. Also, Chen et al. [8] presented an inverse problem in estimating the space dependent thermal conductivity of a functionally graded hollow cylinder. They concluded that the inverse Levenberg-Marquardt method (LMM) are still reliable as the measurement errors are considerable low. Later, Xu et al. [9] investigated length-dependent thermal conductivity in suspended single-layer graphene. They found that at 300K, thermal conductivity keeps increasing and remains logarithmically divergent with sample lengths much larger than average phonon mean free path. In other work, Louis [10] analyzed the effect of space-dependent thermal conductivity on the steady central temperature of a cylinder. Other related articles where thermal and heat transfer analysis in cylindrical geometries are carried out can be found in [11-16]. The aim of this present work is to analyse the impact of variable thermal

conductivity on temperature distributions and rate of heat transfer of some carefully selected fluid which are of high engineering applications. Section one gives the introduction of the work while sections 2, 3 and 4 respectively give the mathematical analysis, results and discussion and conclusion.

Mathematical Analysis

Consider a porous annulus bounded by two vertical cylinders of radius a and b respectively for the inner and outer cylinder. Fluid

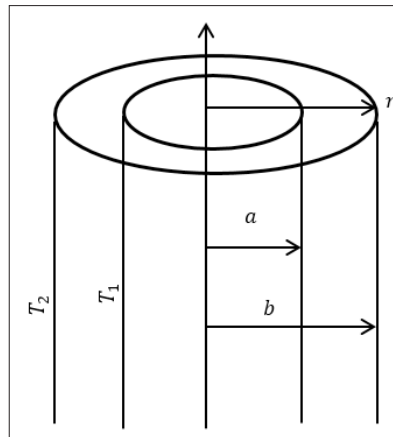


Figure 1: Physical model and coordinate system.

where h is the specific enthalpy, ϕ is the dissipation function, T is the absolute temperature and V is the radial velocity, ρ is the fluid density and k is the thermal conductivity. Neglecting the specific enthalpy and dissipation function and also assuming that the flow is incompressible and along -direction only (thermally fully developed); we have:

$$\rho C_p \left(\frac{\partial T}{\partial t} + V \frac{\partial T}{\partial r} \right) = \frac{1}{r} \left(\frac{\partial}{\partial r} \left[k(r) \left[r \frac{\partial T}{\partial r} \right] \right] \right) \quad (2)$$

In this problem, we assume the temperature distribution to be steady and thermally fully developing with constant radial velocity $r = \frac{a}{r}$. Since the temperature of inner surface of the outer is greater than that of outer surface of inner cylinder $T_2 > T_1$, the thermal conductivity is assumed to vary directly with space by the linear function

$$k(r) = k_0 (1 + b_1 r) \quad (3)$$

where b_1 is the coefficient of variation of the fluid thermal conductivity (1/m). It is well known that the thermal conductivity of gases increases with increase in temperature, and the reverse for liquid. Hence the fluid closer to surface $r=b$ is assumed to have higher thermal conductivity than the fluid closer to the surface

$r=a$.

Under these assumptions, the energy equation reduces to:

$$\frac{V_0 a}{r} \frac{dT}{dr} = \frac{k_0}{\rho C_p r} \left(\frac{d}{dr} \left[(1 + b_1 r) \left(r \frac{dT}{dr} \right) \right] \right) \quad (4)$$

is assumed to be entering the annulus through the surface $r=a$ and going out at the surface $r=b$ at the same rate (continuity satisfied). The inner surface of outer cylinder ($r=b$) is assumed to be heated with temperature (T_2) greater than the surrounding fluids having temperature (T_0) and the outer surface of inner cylinder ($r=a$) is heated to a temperature (T_1) as shown in (Figure 1). The energy equation governing the temperature distribution is given by:

$$\rho \left[\frac{\partial h}{\partial t} + \nabla \cdot (h \vec{V}) \right] = - \frac{D\rho}{Dt} + \nabla \cdot (k \nabla T) + \phi \quad (1)$$

Subject to

$$T = T_1 \text{ at } r = a$$

$$T = T_2 \text{ at } r = b \quad (5)$$

Using the following dimensionless quantities in equations (4),

$$\theta = \frac{T - T_0}{T_2 - T_0}, \quad \lambda = \frac{b}{a}, \quad B = b_1 a, \quad R = \frac{r}{a}, \quad \sigma_r = \frac{T_1 - T_0}{T_2 - T_0}, \quad Pr = \frac{\nu}{\alpha}, \quad S = \frac{V_0 a}{\nu} \quad (6)$$

$$\frac{1}{R} \left(\frac{d}{dR} \left[(1 + BR) \left(R \frac{d\theta}{dR} \right) \right] \right) - \frac{S Pr}{R} \frac{d\theta}{dR} = 0 \quad (7)$$

subject to the following dimensionless boundary conditions

$$\theta = \sigma_r \text{ at } R = 1$$

$$\theta = 1 \text{ at } R = \lambda \quad (8)$$

where Pr is the Prandtl number, S is the suction or injection parameter, T_0 is the ambient temperature, B is the dimensionless number governing the variation of thermal conductivity ($B \geq 0$) and σ_r is cylinder surface temperature difference ratio. It is good to note that when $B=0$, it corresponds to constant thermal conductivity [Jha and Babatunde [16].

Solving Eq. (7) subject to (8), we have:

$$\theta(R) = \frac{1}{((1+B)\lambda)^{SPr}} \left\{ X_1 + X_2 \left(\frac{R}{1+BR} \right)^{SPr} \right\} \quad (9)$$

where

$$X_1 = \sigma_r [(1+B)\lambda]^{SPr} - [1+B\lambda]^{SPr}, \quad X_2 = (1-\sigma_r) [(1+B\lambda)(1+B)]^{SPr} \quad (10)$$

The Nusselt number at the surfaces of the cylinders is given in dimensional form as:

$Nu = \frac{h a}{k}$, where h is the heat transfer coefficient and is defined as

$$h = k \frac{dT}{(T_2 - T_0)} \quad (11)$$

So that by the use of Eq. (6), Eq. (11) becomes

$$\left. \frac{d\theta}{dR} \right|_{R=1} = \frac{S Pr X_2}{(1+B)^{SPr+1}} \quad (12)$$

$$\left. \frac{d\theta}{dR} \right|_{R=\lambda} = \frac{S Pr X_2 \lambda^{(Spr-1)}}{(1+B)^{SPr+1}} \quad (13)$$

Results and Discussion

The solutions obtained are seen to be govern by Prandtl number (Pr), thermal conductivity variation parameter (B), suction/injection parameter (S), radius ratio (λ) and cylinder surface temperature difference ratio (σ_r). We have considered different applicable fluids dictated by Prandtl number: Liquid Lithium are used to cool the fusion reactors blankets, it is also used for superconducting magnets, such as those needed in nuclear magnetic resonance imaging and nuclear magnetic resonance Liquid Lithium (Pr=0.065), Helium (Pr=0.261), the boiling point of helium is the least when compared to any other liquid. It is used to obtain the lowest temperatures required in lasers. Helium is also used in nuclear reactors as a cooling gas and used as a flow-gas in liquid-gas chromatography. It finds its application in airships and helium balloons. In addition, helium balloons are used to check the weather of a particular region. Helium is preferred over hydrogen though hydrogen is cheaper, as helium is readily available and hydrogen is highly inflammable. Hence due to safety issues helium is preferred in aircrafts [17]. Oxygen (Pr=0.63), air (Pr=0.71) and ammonia (Pr= 1.38) on the other hand is use as cleaning, bleaching, deodorizer, drugs, dyes and fertilizer Ried et al. [17].

Figure 2 depicts temperature profile at different values of thermal conductivity variation parameter (B) for different fluids at fixed value of suction/injection and radius ratio. It is seen that temperature is highest in Liquid Lithium (which has the lowest Prandtl number) and lowest for water (which has the highest Prandtl number). This can be attributed to the fact that increase in Prandtl number always reduce fluid temperature. In addition, the role of space dependent thermal conductivity is to increase fluid temperature in the annulus. Figure 3 exhibits combined effect of cylinder temperature difference ratio (σ_r) and thermal conductivity variation parameter on fluid temperature. It is observed that fluid temperature increases with increase in (σ_r). Moreover, for higher value of (σ_r), (which almost corresponds to symmetric heating), the significant of is seen to be reduced compare of lower value of (σ_r) ($\sigma_r \leq 0.5$). Figures 4 & 5 present the impact σ_r and B for different fluids at the outer surface of inner cylinder and inner surface of cylinder, respectively. From both figures, the rate of heat transfer is seen to increase with increase in B but decreases with increase in σ_r . In

addition, the rate of heat transfer is highest for atmospheric air and lowest for liquid Lithium. This explains why Liquid Lithium and helium are used as nuclear plant coolant.

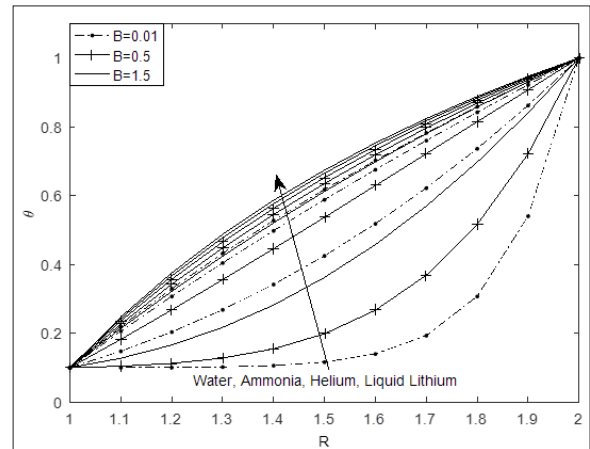


Figure 2: Temperature profile for different fluids varying B at S=2.0, $\sigma_T=0.1$, $\lambda=2.0$.

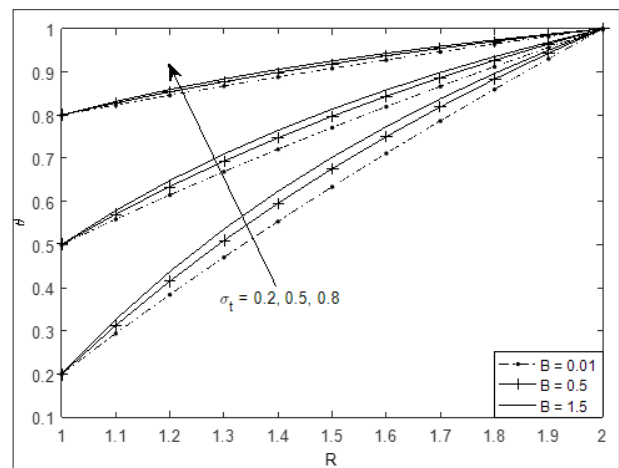


Figure 3: Temperature profile of liquid lithium varying B and σ_T at $\lambda=2.0$, S=2.0.

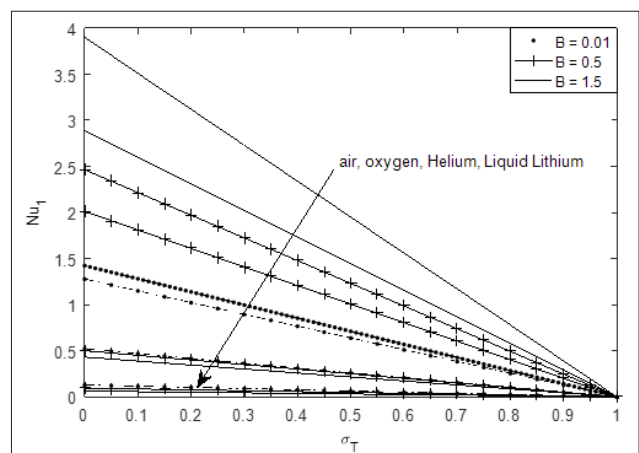


Figure 4: Nusselt number of fluids for different values of B and σ_T at S=2.0, $\lambda=2.0$, R=1.

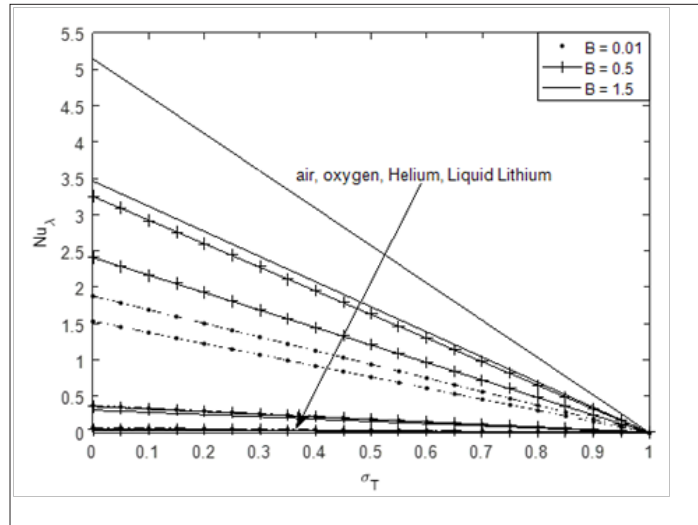


Figure 5: Nusselt number of fluids for different values of σ_T and B at $S=2.0, \lambda=2.0, R=2.0$.

Table 1: Rate of heat transfer of different fluid for different values of σ_T and B at $S=2.0, \lambda=2.0$.

Fluid	σ_T	B	Nu_1	Nu_λ
Liquid Lithium	0.2	0.01	0.1032	0.0562
		0.5	0.0759	0.0415
		1.5	0.0498	0.0273
	0.5	0.01	0.0645	0.0353
		0.5	0.0474	0.0259
		1.5	0.0311	0.017
	0.8	0.01	0.0258	0.0141
		0.5	0.019	0.0104
		1.5	0.0125	0.0068
Helium	0.2	0.01	0.4176	0.2998
		0.5	0.3998	0.287
		1.5	0.3444	0.2473
	0.5	0.01	0.261	0.1874
		0.5	0.2499	0.1794
		1.5	0.2153	0.1546
	0.8	0.01	0.1044	0.075
		0.5	0.0999	0.0718
		1.5	0.0861	0.0618
Oxygen	0.2	0.01	1.0095	1.2089
		0.5	1.6094	1.9272
		1.5	2.3127	2.7694
	0.5	0.01	0.631	0.7556
		0.5	1.0059	1.2045
		1.5	1.4454	1.7309
	0.8	0.01	0.2524	0.3022
		0.5	0.4024	0.4818
		1.5	0.5782	0.6932

Ammonia	0.2	0.01	2.218	7.5123
		0.5	9.9713	33.7725
		1.5	40.527	137.2642
	0.5	0.01	1.4431	4.4877
		0.5	6.232	21.1078
		1.5	25.3294	85.7902
	0.8	0.01	0.5772	1.9551
		0.5	2.4928	8.4431
		1.5	10.1318	34.3161

Table 2: Validation of fluid temperature result using liquid lithium at different R and σ_T and B at S=2.0, $\lambda=2.0$.

σ_T	R	Temperature (Constant Thermal Conductivity)	Temperature (Present Work) B \rightarrow 0
0.2	0.2	0.3832	0.3832
	0.4	0.5524	0.5524
	0.6	0.7103	0.7103
	0.8	0.859	0.859
0.5	0.2	0.6145	0.6145
	0.4	0.7202	0.7202
	0.6	0.8189	0.8189
	0.8	0.9119	0.9119
0.8	0.2	0.8458	0.8458
	0.4	0.8881	0.8881
	0.6	0.9276	0.9276
	0.8	0.9647	0.9647

Table 1 gives the values for rate of heat transfer for different values of governing parameters. It is observed that Nusselt number of fluid with lower Prandtl number (Lithium and Helium) decreases with increase in σ_T and B. The reverse trend is noticed for fluid with higher Prandtl number. (Table 2) on the other hand presents the validation of the present work with the absence of variable thermal conductivity $B \rightarrow 0$. This comparison gives an excellent agreement.

Conclusion

This paper examines an exact solution to show the impact of space dependent thermal conductivity on some fluids in a porous annulus with asymmetric surface heating. The governing energy equation is derived and closed form expressions are obtained for temperature and rate of heat transfer. Based on depicted graphs and tables, the following major conclusions can be drawn:

- Fluid temperature as well as rate of heat transfer increase with increase in thermal conductivity variation parameter.
- Temperature is higher for fluids with lower Prandtl number (liquid lithium) and lower for fluids with higher Prandtl number (water)

c) Fluid temperature as well as rate of heat transfer becomes independent of governing parameters for higher value of σ_T , ($\sigma_T \approx 1$).

d) It is hoped that the results obtained in this paper will not only provide useful information for industrial applications (quality control during manufacturing process), but also serve as an improvement on the previous studies.

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