

Room Temperature Photoluminescence and Optical Parameters of Micro Rods ZnO Thin Films Deposited on Silicon Substrates

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Abstract

Micro structured zinc oxide (ZnO) is of potential use in solar cells, solid state gas sensors, and other optoelectronic devices. For the advancement of these technologies, a comprehensive understanding of these materials' optical characteristics is considered essential. Using the spray pyrolysis (SP) method, thin films of ZnO micro rods are deposited from aqueous zinc chloride solution on silicon substrates held at 350 ± 5 °C. The films have been studied using the data acquired from the X-ray diffraction (XRD), scanning electron microscopy (SEM), X-ray energy dispersive spectroscopy (EDS), reflectance and photoluminescence (PL) spectroscopy. Hence, the XRD and SEM confirmed the films' hexagonal phase and revealed the tapered ends of the hexagonal rods. The optical parameters including the refractive index, extinction coefficient, absorption coefficient, real and imaginary dielectric constants, conductivity, and the energy loss are all deduced from the reflectance spectra. Finally, the photoluminescence spectra of the studied samples are measured at room temperature using two excitation wavelengths: 300 and 365nm. As such, the obtained results confirmed the quality of the prepared films and allowed proper determination of their respective optical characteristics.

Keywords: ZnO micro rods; Optical parameters; Spray pyrolysis; Solar cells; Photoluminescence

Introduction

Zinc oxide is a multipurpose material, which assumes the Wurtzite structure with space group $C6v4$ or $P63mc$ [1,2]. It is redeemed useful in so many applications including solar cells, solid state gas sensors, photovoltaic and other optoelectronic devices, and in the biomedical industry as well [3-6].

Although numerous preparation techniques including sol-gel [7], electrodeposition [8], hydrothermal [9], chemical bath deposition (CBD) [10], solvothermal synthesis [11], RF sputtering [12], glaze firing [13], reactive sputtering [14], and many others are considered adequate to produce ZnO in its intended thin film forms, the spray pyrolysis in particular has received recently adequate attention because of its simplicity, economics, and ease of scale up for large scale production [15-20]. In addition, the spray pyrolysis method can be easily altered to produce films of different morphologies including nano-sheets, -wires, -combs, -flowers, and -rods [6]. This happens as a direct consequence of the preparation methods and conditions and usually affected by the preferential growth along the c-direction [1,6].

Even though ZnO occurs in nature as an n-type with large exciton binding energy of 60meV and wide bandgap energy of 3.37eV at room temperature [21], it can nowadays be easily synthesized in either n- or p-type by intentionally doping it with elements of group-III or group-I and V, respectively. Nevertheless, its brilliant electrical conductivity is a direct consequence of the zinc excess at the interstitial positions and/or due to the oxygen vacancies [17]. In a previous work, the authors studied the functionalities, compositional quality, and the chemical bonding of the ZnO thin films using FTIR spectroscopy. They observed the Zn-O bond stretching vibration peaks at 420cm^{-1} , 454cm^{-1} , 485cm^{-1} , and 546cm^{-1} , and the substrate's Si-O stretching vibration at around 1072cm^{-1} as well [16].

In this work, the authors extend their previous studies to explore in more details both the optical parameters of the prepared ZnO thin films and their room temperature

photoluminescence spectra. Both the optical parameters, as deduced from the reflectance spectra, and the measured PL are considered major contributors to the quality of the prepared films and help in providing proper understanding of the films' optical characteristics that are needed to improve their usage in the optoelectronic industries.

Experimental Procedure

Thin films of ZnO were prepared, using the spray pyrolysis method, on silicon substrates at a substrate temperature of 350 ± 5 °C. The single-side polished, boron-doped p-type, 3-30 Ω cm with (001) orientation, $675 \pm 25 \mu\text{m}$ thick Si substrates of area of $5 \times 5 \text{mm}^2$ were cut from a 150mm wafer (Semiconductor Wafer, Inc.). Prior to heating the substrates in preparation for the deposition process, the wafers were washed with acetone, methanol, DI water, and blown dry. The full preparation protocol is published elsewhere [16]. X-ray diffraction patterns were measured utilizing a Shimadzu diffractometer (Model 7000). The thin films were examined by an FEI scanning electron microscope (Inspect F50) equipped with a Bruker EDS attachment. The room temperature optical reflectance is measured at normal incidence in the wavelength range 240-840nm using the FilmTek-3000 characterization tool. The measured data is later fitted using various models while the films' thicknesses are given as byproducts of the fits. Hence, the thicknesses of films 1 and 2 are found \sim (200 and 120) nm, respectively. Subsequently, each film reflectance was used to deduce the thin film's optical parameters including the refractive index, extinction coefficient, absorption coefficient, real and imaginary dielectric constants, conductivity, and the energy loss. Finally, PL spectra at room temperature were recorded with a Cary Eclipse (Varian) Spectrofluorometer under excitation wavelengths of 300 and 365nm. While the spectral monochromator emission slit was kept at 5nm, the excitation slit was opened at 5, 10, or 20nm. For the excitation wavelength 300nm, the PL data was recorded within the range 330-570nm, while for the excitation wavelength 365nm, the PL data was recorded at 390-670nm.

Results and Discussion

The X-ray patterns of the thin films of ZnO deposited by the spray pyrolysis agree well with the reference pattern in the Joint Committee on Powder Diffraction Standards data base (JCPDS card 00-036-1451) as no extraneous diffraction peaks appeared, indicating a phase-pure Wurtzite structure [16]. Calculation of unit cell parameters from the diffraction pattern gave values of $a_0 = (3.2657 \text{ and } 3.2659 \pm 0.0061) \text{ \AA}$ and $c_0 = (5.213 \text{ and } 5.210 \pm 0.057) \text{ \AA}$ compared with the values of $a_0 = 3.249 \text{ \AA}$ and $c_0 = 5.206 \text{ \AA}$ reported by McMurdie et al. [22] SEM images confirmed the hexagonal microrods nature of the well faceted crystals, in a range of particle sizes of much less than $1 \mu\text{m}$ and the Films' composition from EDS-analysis found to consist of Zn and O mainly in addition to small concentration of Cl [16].

From the reflectance (R) of the two ZnO thin films deposited on silicon substrates as a function of radiation wavelength, measured

and published elsewhere [16], it appeared that R decreases as the wavelength of the incident radiation increases with two pronounced maxima. These maxima are the same as observed in the reflectance of a bare Si native oxide. The appearance of these two maxima is due to the very thin layer ($\sim 10 - 15 \text{ \AA}$) of the native silicon oxide formed on the top of the silicon substrates (SiO_x : $x \approx 1.5$). However, the direct influence of the ZnO layers on the overall reflectance of the samples is obvious in lowering the reflectance values based on the different thicknesses of the two samples.

The transmittance T was calculated from the reflectance using the relation: $T+R=1$, then used to find the absorption coefficient α , where

$$\alpha = -\ln T / t \quad (1)$$

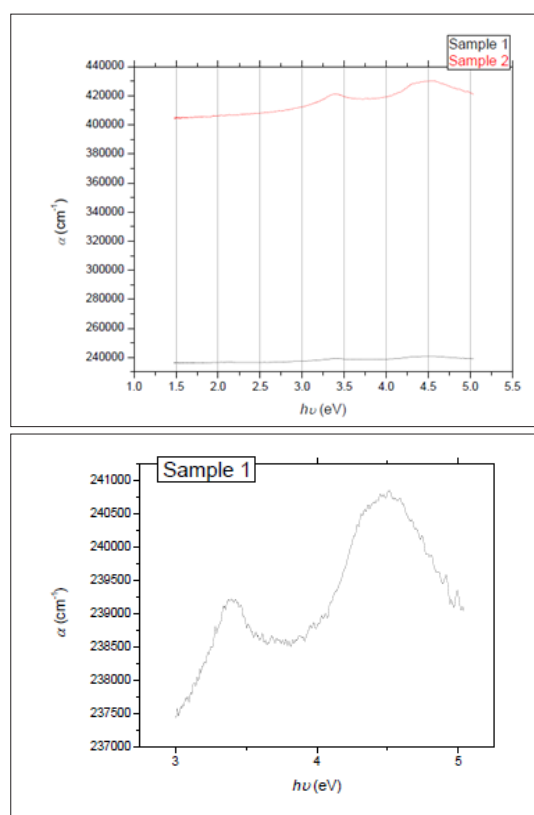


Figure 1: (a) Alpha of the two ZnO thin films deposited on Si substrates. (b) Exploded view of Sample 1.

Where t representing the thickness of each film. (Figure 1a & 1b) depict the relationship between α and the photon's energy $h\nu$ (eV). As the figures show, there is a slight increase of α with $h\nu$, however the absorption in sample 1 is much greater compared to sample 2 as it is the thicker one. The two maxima related to SiO_x are apparent in the curve of sample 2 while they are very weak in sample 1.

To find the refractive index and extinction coefficient of the studied samples, the following relationships are utilized for a light wave incident normally from the air on the films 19.:

$$R = \frac{(n^* - 1)^2}{(n^* + 1)^2} = \frac{(n - 1)^2 + k^2}{(n + 1)^2 + k^2} \quad (2)$$

With $n^* = n + i k$, where n is the refractive index and k is the extinction coefficient of the film which also can be calculated using the following relationship based on the absorption coefficient α and the radiation wavelength λ in free space:

$$k = \frac{\lambda \alpha}{4\pi} \quad (3)$$

Moreover, solving equation (2) for the refractive index gives:

$$n = \frac{(1 + R) + (1 + R)^2 - (1 - R)^2 (1 + k^2)^{1/2}}{1 - R} \quad (4)$$

Calculated n and k as a functions of photon's energy $h\nu$ and based on the previous equations were shown in Figure 2 & 3, respectively. Refractive index n , as seen in Figure 2 increases with photons energy, and it is larger for sample 2, which has the smaller thickness. Also, the maxima characteristics of bare Si native oxide are apparent in the curves. Extinction coefficient k decreases with $h\nu$ and reaches about 0.4 for sample 1 (the thicker), and 0.9 for sample 2 (the thinner). The non-zero value of k , is due to impurities in the samples. It is observed that k is larger for sample 1, where the larger thickness results in more absorption.

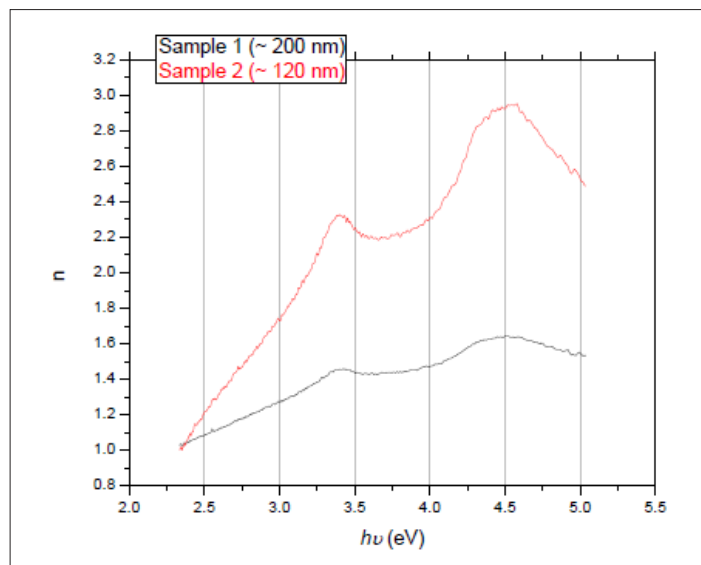


Figure 2: Refractive indices of the two ZnO thin films deposited on Si substrates.

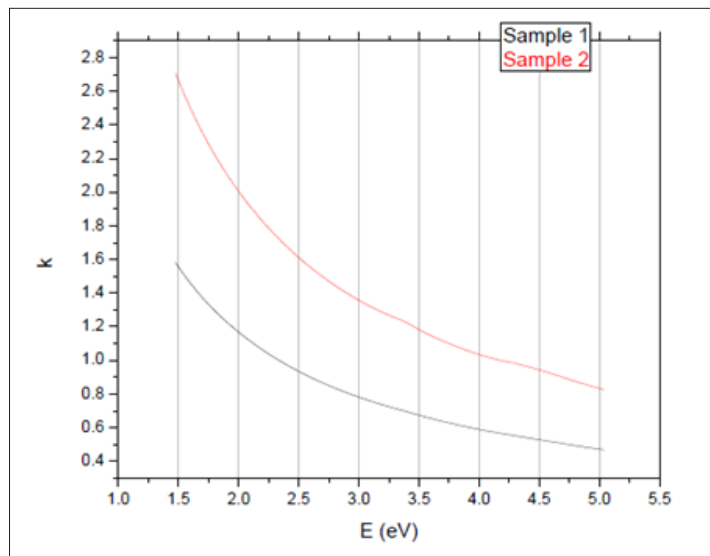


Figure 3: Extinction coefficient of the two ZnO thin films deposited on Si substrates.

Polarizabilities of solids are related to their dielectric constants where any interaction between the photons and the electrons within the films are reasonably reflected by the variation of the dielectric constants with the photons' energies. Both real and imaginary components of the complex dielectric constant can be expressed conveniently as given in equations 4 and 5, respectively:

$$\epsilon_1 = n^2 - k^2 \quad (5)$$

And

$$\epsilon_2 = 2nk \quad (6)$$

Figure 4 displays how ϵ_1 varies with the photon's energy $h\nu$ for ZnO. As shown in the Figure, ϵ_1 increases slowly with $h\nu$ with two

obvious humps. This abrupt increase indicates strong interaction of photons and electrons with maxima at these two energies. Figure 5 displays the imaginary part of the dielectric constant ϵ_2 as a function of the photon's energy $h\nu$ for the ZnO film. It is noticeable that ϵ_2 behaves exactly like k as it reflects the free carriers' absorption. The two maxima in ϵ_2 occur at the same photons' energy as in ϵ_1 . Heating the dielectric material in a varying electric field results in an energy loss which is normally represented by the following relationship:

$$\tan\delta = \epsilon_2 / \epsilon_1 \quad (7)$$

and the loss angle is given as:

$$\delta = \tan^{-1}(\epsilon_2 / \epsilon_1) \quad (8)$$

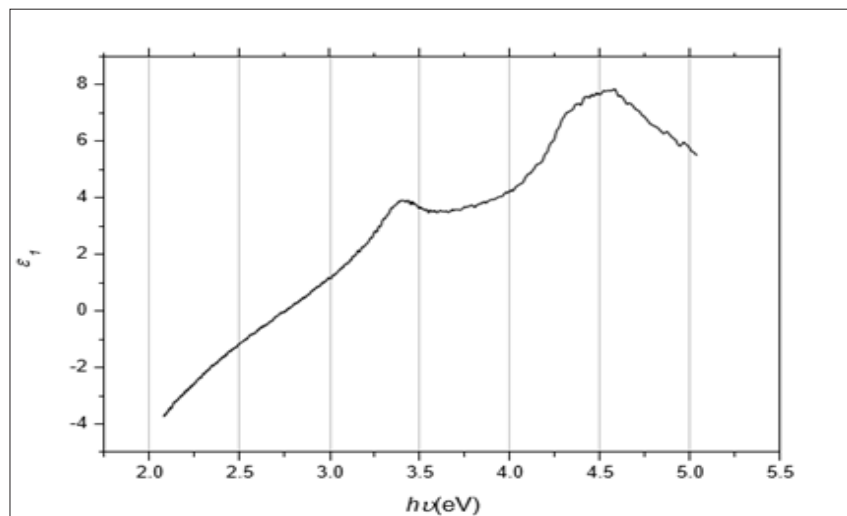


Figure 4: The real part of the dielectric constant ϵ_1 against the photon's energy $h\nu$ for the ZnO thin film deposited on Si substrate.

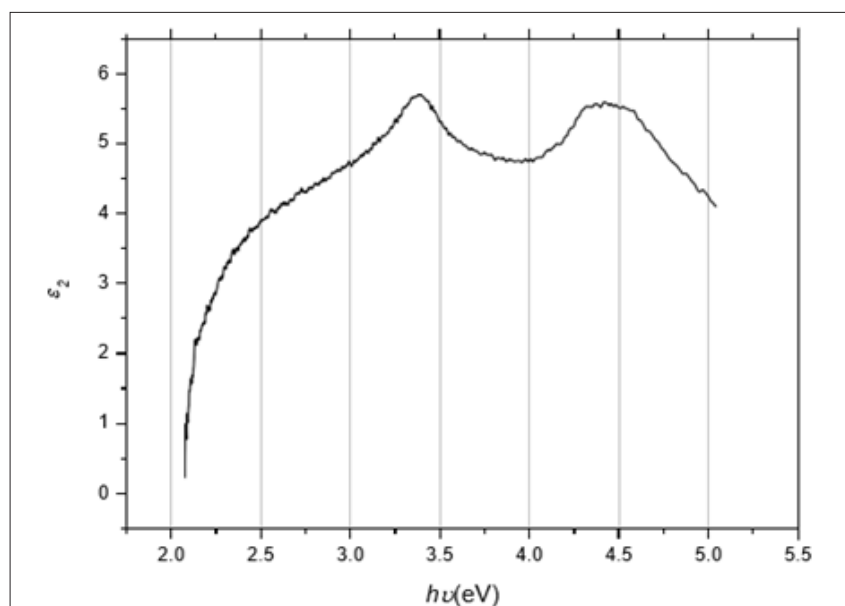


Figure 5: The imaginary part of the dielectric constant ϵ_2 against the photon's energy $h\nu$ for the ZnO thin film deposited on Si substrate.

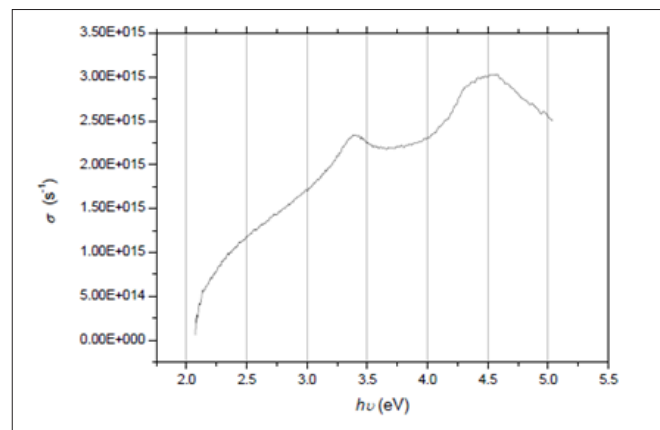


Figure 6: The energy loss as a function of the photon's energy $h\nu$ for the ZnO thin film deposited on Si substrate.

Figure 6 illustrates the relationship between $\tan \delta$ and $h\nu$ for the ZnO film. From the figure, it is shown that the dielectric loss possesses one maximum. The optical conductivity σ of the ZnO thin films can be determined using the formula:

$$\sigma = \alpha nc / 4\pi \quad (9)$$

where c represents the speed of light. Projecting the previous formula on the acquired data resulted in obvious variations in the optical conductivity as a function of the photon's energy $h\nu$ and reveals increasing optical conductivity at two maxima along the high energy side that can be attributed to the films' high absorbance within these regions. Finally, as indicated in all figures, ZnO films of varying thicknesses prepared in this investigation have optical parameters with same features however with different values.

Room temperature characteristic photoluminescence of the deposited films were shown in Figure 7. As shown in Figure 7a & 7b, two main features were seen in the photoluminescence spectra; featured band in the UV region characteristic of the ZnO,

and another emission feature extending in the visible region and originating from all sorts of defects and impurities [1,22]. Opening the slit width resulted in higher registered relative intensity as more energy is allowed to excite the luminescent centers within the matrix. Attempts to explain the PL spectra suggests that the UV luminescence is originated from near band edge transitions. The presence of several peaks within the UV region is considered a confirmation indicator of the presence of different nano particle sizes within the films that experience ZnO nanocrystalline excitonic emission (EE) with quantum confinement effects. Moreover, extrinsic impurities as well as deep levels and defects in ZnO crystal structure (i.e. oxygen and zinc interstitials and oxygen and zinc vacancies) are considered mainly responsible for the visible part of the emission spectrum. The zinc interstitials and zinc vacancies (Zni ,VZn) are related to blue emission. The singly ionized oxygen vacancies are commonly attributed to initiate the green emission as it is easily influenced by surface modification. Both orange and green photoluminescence are related to the existing amount of oxygen in the films not to the dopants' trace amounts [1,4,16].

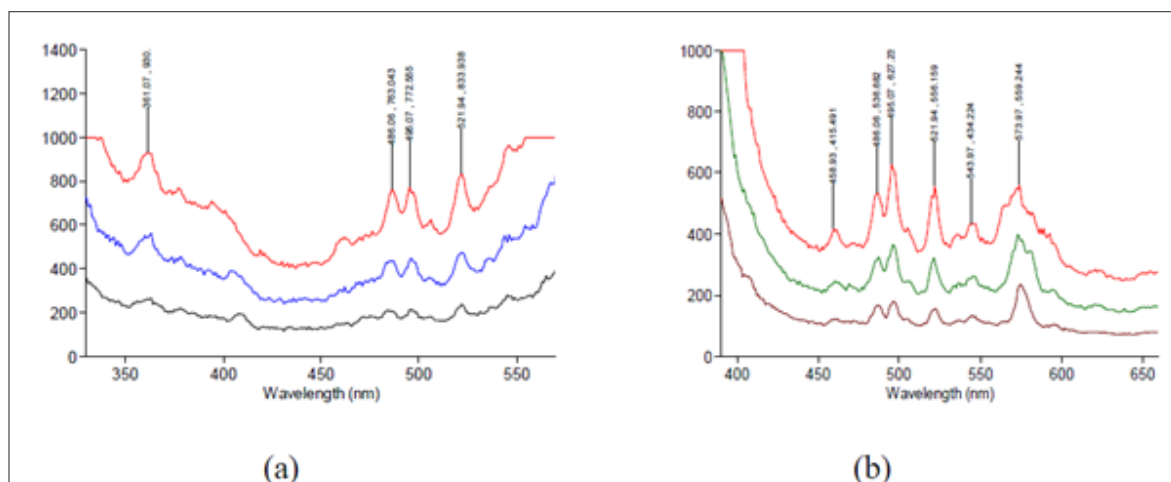


Figure 7: RT-PL spectra of ZnO thin films for excitation wavelength (a) 365nm. and (b) 300nm. Ex slit width = 5, 10, and 20nm, Em Slit width = 5nm.

Conclusion

Thin films of ZnO were produced on silicon substrates held at 350 ± 5 °C. XRD diffraction revealed the hexagonal structure of the films. SEM microscopy confirmed the hexagonal phase and revealed the morphology of the films, where they showed hexagonal micro rods. EDS, reflectance, and PL measurements were collectively utilized to characterize the films. Optical parameters including α , n , k , e_1 , e_2 , σ and the energy loss are all determined based on the reflectance spectra. Measured Room temperature PL spectra presented two featured bands in the UV and visible regions. It is found that varying the thickness of the samples and the presence of impurities have significant effect on the results. The widely opened excitation slits produced the PL of maximum intensity. The measured PL signals consist of UV and VIS luminescence. This study is of great interest to photovoltaics, gas sensors, and optoelectronics industries.

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