

# Synthesis and Magneto-optical Properties of Co doped TiO<sub>2</sub> Nanotubes from Electro-spun Fiber Templates

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## Abstract

Submicron Co doped TiO<sub>2</sub> nanotubes were synthesized by the tubes by fiber template approach. Polymer fiber templates were fabricated by electrospinning while the tubes were synthesized by sol-gel deposition followed by thermal degradation of the polymer core. The diameters of the tubes ranged between ~350 ± 100nm with an average wall thickness of ~100 ± 50nm from SEM and TEM analysis. PXRD analysis indicated Co was homogeneously distributed within substitutional sites of the polycrystalline anatase phase of the TiO<sub>2</sub>. Photoluminescence studies confirmed n-type doping as indicated by reduction in intensity of highest energy direct photoemission (~420nm) due to charge transfer as well as presence of photoemission at (~470nm) associated with the presence of oxygen vacancies. Magnetic studies results indicated that the Co doped TiO<sub>2</sub> tubes were paramagnetic at room temperature.

## Introduction

Titanium dioxide, a wide band gap semiconductor, has stimulated intense research interest due to its excellent optical transmission in the visible and near infrared region, high refractive index, high dielectric constant and useful photo catalytic properties [1]. As a result, it has found many interesting device applications such as photovoltaics, optics, magnetic storage, catalysis, sensors and battery electrodes [2-4]. More recently, Co doped titanium oxide has stimulated intense and broad studies both in theory and experimental as a dilute magnetic semiconductor following discovery of robust, room temperature ferromagnetism important in spintronics [5,6].

The fabrication of dilute magnetic semiconductors has been achieved using several deposition techniques including oxygen assisted molecular beam epitaxy, pulsed laser deposition, solid vapor deposition/co-sputtering, ion implantation and cathodic electrolytic deposition [5,7-9]. However, these methods are not only expensive but less straightforward hence have limitations in terms of throughput. Similarly, the origin of room temperature ferromagnetism still remains controversial in dilute magnetic semiconductors prepared by these techniques in reducing environments or ultra-high vacuum because of the possibility of segregation of transition metal particles resulting in phase separation [10].

Currently, endeavour's particularly focused on wet chemical synthesis under suitable conditions have proved quintessential in addressing these limitations while aiding in understanding the origin of ferromagnetism. Consequently, not only a single phase of Co doped TiO<sub>2</sub> polycrystalline are produced but also allows for integration with Si technology that is important in spintronics [10,11]. In this work Co doped TiO<sub>2</sub> nanotubes were prepared by the

tubes by fiber templating technique via a Sol-gel process. Electrospun fibers served as templates which on thermal degradation yielded hollow Co-TiO<sub>2</sub> nanotubes [12,13]. Electrospinning is a technique that relies on repulsive electrostatic forces to produce solid fibers in the nanometre to micro-meter range from polymer solutions or melts [14]. The forces offer unique advantages for preparing homogenous multicomponent oxides compared to gas phase techniques [15].

The Co doped TiO<sub>2</sub> were characterized for magnetic and optical properties relative to thin films obtained by molecular beam epitaxy and co-sputtering techniques. Similarly, magneto-optical properties of Co doped TiO<sub>2</sub> nanotubes were of importance since most studies reported have focused on 2 dimensional materials, thin films. The tubular structure of these Co doped TiO<sub>2</sub> nanotubes becomes particularly important due to its high aspect ratio, high porosity and the quantum confinement effects as a result of reduction in dimensions.

## Experimental Section

### Materials

Poly lactide (PLA) pellets (Mwt. 180,000), Polycarbonate pellets (Mwt. 64,000) methylene chloride, N, N-dimethylformamide (DMF), Tin Chloride Palladium Chloride, Cobalt (II) nitrate hexahydrate, Titanium isopropoxide (TiP), 2-propanol and Hydrochloric acid all from Sigma-Aldrich were used as received.

### Electrospinning and Sol-gel deposition

The template fibers were prepared by electrospinning poly lactide or polycarbonate polymer solution. The polymer solutions were prepared by dissolving 180mg/mL in dichloromethane/dimethylformamide (0.65/ 0.35) solvent mixture. Electrospinning was done at 20kV (applied voltage) with a working distance of 20cm between the collector screen and the spinneret [14]. The electro-spun fibers were peeled off from the aluminium foil collector screen after soaking in 1M HCl solution and then rinsed in deionized water prior to use.

Titanium oxide poly lactide coaxial fibers were fabricated via wet chemical synthesis through the hydrolysis of titanium isopropoxide, TiO<sub>2</sub> precursor, followed by condensation in the sol-gel deposition process. The template was first sensitized and then activated by immersion in 3.0mM PdCl<sub>2</sub> aqueous solution containing 0.01M HCl. The colloidal suspension for Co doped TiO<sub>2</sub> nanotubes was prepared using Titanium isopropoxide (TiP) as TiO<sub>2</sub> precursor and Cobalt (II) Nitrate hexahydrate as the dopant source. Under constant stirring, 1.0mL of TiP was added drop wise to a 20mL solution of 0.01M (0.055g) Cobalt (II) Nitrate hexahydrate in 2-propanol. The resulting mixture was stirred for 2h over an ice bath to form a homogenous and stable colloidal solution. The percent dopant loading was varied by varying the concentration of Co(NO<sub>3</sub>)<sub>2</sub>·6H<sub>2</sub>O used.

Activated and sensitized template fibers were immersed into the colloidal suspension and the reaction proceeded with constant stirring for 2h over an ice bath resulting in a coating of Co-TiO<sub>2</sub>. After coating, the fibers were removed from the colloidal suspension

and allowed to hydrolyze in air of 24h at room temperature. This resulted in gelation and formation of PLA-Co-TiO<sub>2</sub> coaxial fibers. The coaxial fibers were then calcined in air at 500 °C at a ramp rate 10 °C/min and annealed at that temperature for 3h.

### Instrumentation

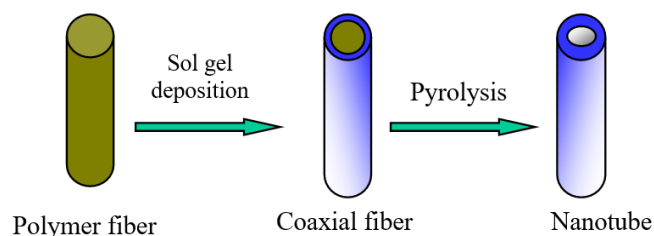
The morphology, dimensions and elemental composition of the synthesized materials were characterized by electron microscopy. Scanning Electron Microscope (SEM) model Hitachi S-570LB equipped with an Energy Dispersive Spectroscopy (EDS) was used. For morphological studies lower acceleration voltage of 5kV was applied while for EDS 15-20kV was used. For non-conducting samples a thin film of Au/Pd was coated on the surface prior to taking image. Transmission electron microscopy (TEM) was performed on Hitachi 2000 TEM instrument. Thermogravimetric Analysis (TGA) was performed on TA 2950 Analyzer.

The electronic and crystal structure was analyzed using spectroscopic techniques. Electronic absorption spectra were recorded on a Hewlett Packard 8452A UV-visible spectrophotometer. Infra-red spectra were recorded on a digilab FTS-40 PRO as a KBr pellet or fiber mat. The X-ray diffraction spectra were measured on finely ground tubes using Scintag X-ray diffractometer with X-ray wavelength  $\lambda$  1.5418 Å (Cu K $\alpha$ ) radiation source. The samples for PXRD were smoothly ground and compacted to at least 1mm in thickness to prevent penetration of the X-ray beam. The crystalline size was estimated by applying the Scherrer equation to the FWHM of the (101) peak of anatase.

Magnetic properties were determined on a Superconducting Quantum Interference Device (SQUID) magnetometer, Quantum Design MPMS XL-5. The magnetic susceptibility ( $\chi=M/H$ , M is magnetization, H is applied magnetic field) of the samples was measured from 2 to 400 K in a magnetic field of 1000 Oe. Magnetization curves were measured at 298K in magnetic fields up to 5T. The samples were zero-field-cooled to 5K before the magnetization measurements.

## Results and Discussion

### Tubes by fiber templating and material characterization



**Scheme 1:** Schematic diagram of tubes by fiber template process

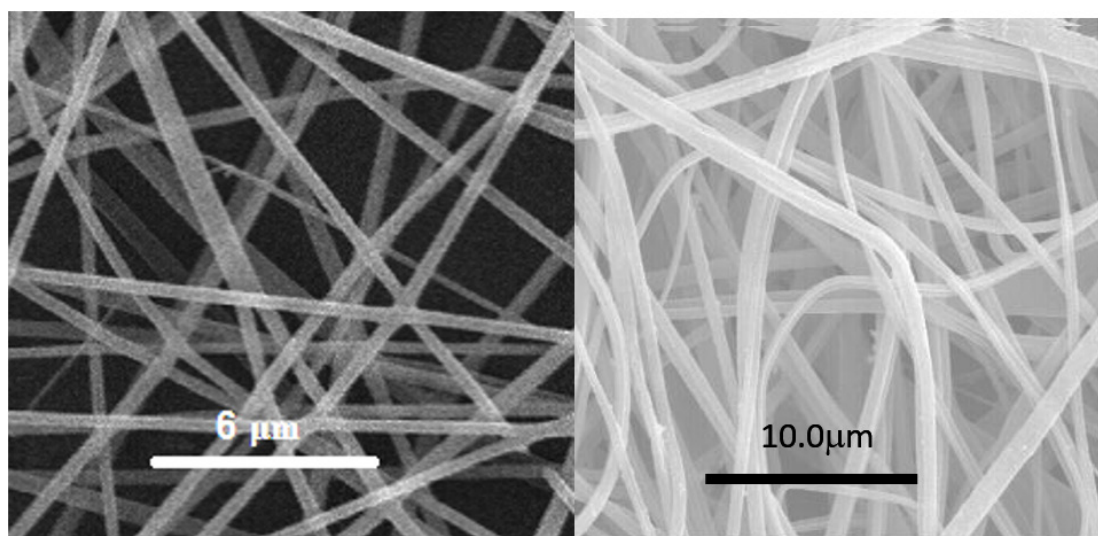
Templating remains one of the most rapidly growing areas of research in materials synthesis. It entails fabrication of structured materials over a scaffold material, template, such as electro-spun fibers. The removal of the template by thermal treatment generates the desired structure in this case submicron tubes as shown in scheme 1 [12,16,17]. The template-based approach provides

a versatile and a low-cost method of preparing nanometer to submicron materials in high volume and with better control of dimension and surface morphology from nanometer to submicron length scale templates.

Sol-gel technology is an attractive methodology for the design and synthesis of advanced ceramic material with anisotropic properties through the control of both the structure and chemistry of the materials. In particular, doping metal oxides with transition metal ions modifies crystalline structure which determines the physical and chemical properties of these materials [10,11]. Sol-gel technology offers the advantage of low cost, low processing

temperatures, and potential for highly homogenous and quality deposits [17].

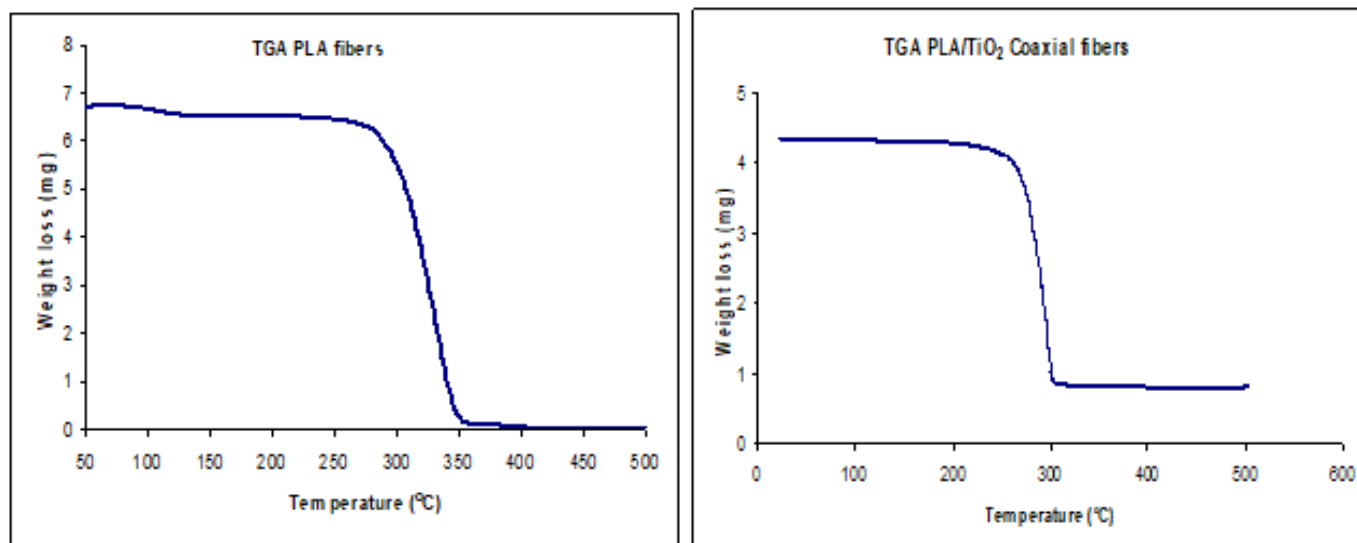
The electrospun PLA or PC nanofibers were preferred as templates materials with  $\text{CH}_2\text{Cl}_2$  and DMF as solvents since they both have a good solubility in methylene dichloride ( $\text{CH}_2\text{Cl}_2$ ). DMF was used as a co-solvent to increase jet stability against decomposition under a high electric field since it has a high dielectric constant. The morphology of PLA fibers was examined using Scanning Electron Microscopy (SEM), revealed that the surfaces were smooth and uniform without beads with an average diameter around  $250 \pm 100 \text{ nm}$ , Figure 1.



**Figure 1:** SEM image of electrospun PLA template fibers (a) Polycarbonate fibers (b).

PLA was selected as template material based on its favorable stability and relatively low decomposition temperature 235-255 °C as confirmed by Thermogravimetric Analysis [18] Figure(2a,2b). Low decomposition temperature was important in reducing cross linking of titanium oxide tubes during thermal treatment. From the

TGA curve of PLA/ $\text{TiO}_2$  coaxial fibers, the initial weight loss around 100 °C to 260 °C was attributed to loss of solvent, decomposition of titanium isopropoxide and decomposition of polylactide. There was no further weight loss after 300 °C indicating the formation of  $\text{TiO}_2$  tubes.

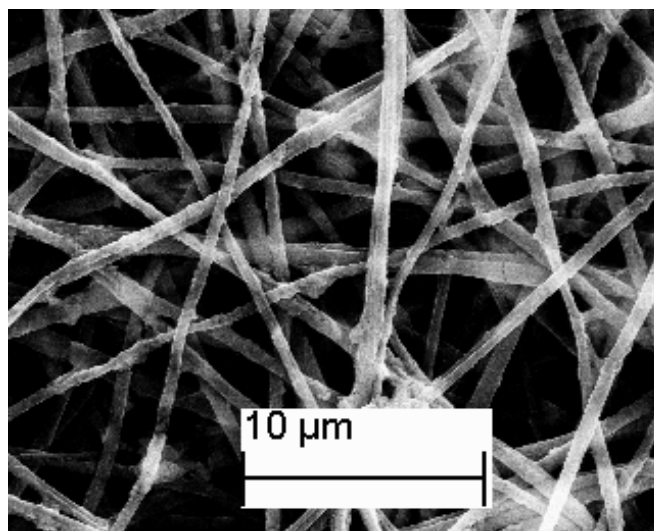


**Figure 2:** Thermogravimetric Analysis of PLA fibers (a) PLA/ $\text{TiO}_2$  coaxial fibers (b).

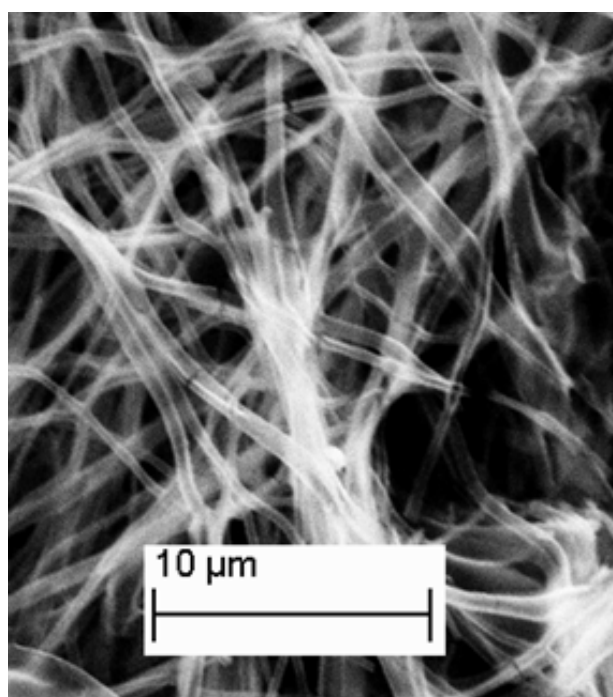
The sol-gel deposition procedure employed was simpler than established gas phase methods, since the sols consisted of only three components: titanium isopropoxide precursor, Cobalt (II) nitrate hexahydrate as the dopant, and 2-propanol as solvent. Since titanium isopropoxide and  $\text{Co}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$  were highly soluble in 2-propanol hence no heating was necessary for the starting mixture. Heating the sol precursors increases the possibility of inadvertently altering the oxidation state of the dopants or forming dopant oxide secondary phases prior to  $\text{TiO}_2$  deposition, as confirmed by PXRD. In addition, no additive was required to improve sol stability and homogeneity. This favored an even drying and uniaxial crystallization of  $\text{TiO}_2$  from SEM characterization.

The resultant PLA-Co- $\text{TiO}_2$  coaxial fibers obtained were found

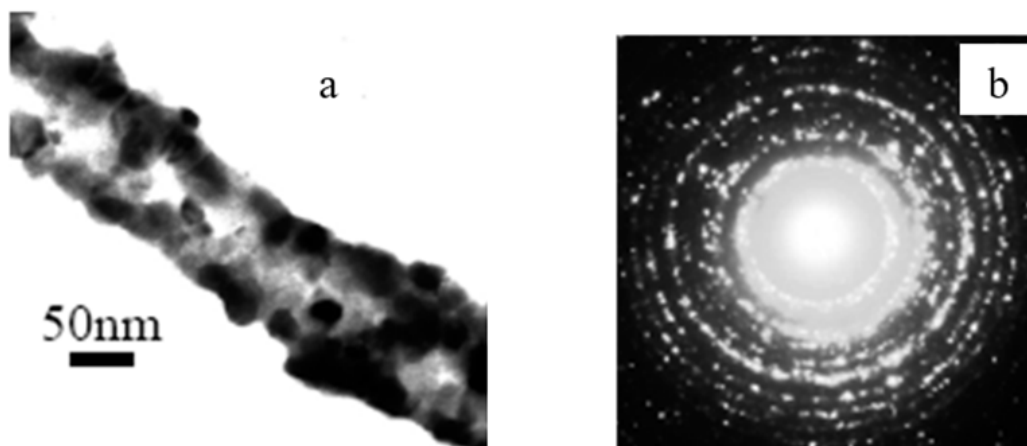
to have a larger diameter compared to PLA template fibers from electron microscopy studies. Figure 3 shows the SEM of PLA-Co- $\text{TiO}_2$  coaxial fibers with diameters ranging between  $450 \pm 100\text{nm}$ . Thermal degradation of PLA template core yielded hollow Co- $\text{TiO}_2$  nanotubes with diameters ranging between  $350 \pm 100\text{nm}$ , Figure 4. The wall thickness of the tubes ranged between  $100 \pm 50\text{nm}$ . The decrease in tube diameter relative to coaxial fibers was attributed to crystallization of  $\text{TiO}_2$ . TEM images and electron diffraction pattern confirmed the tubes were hollow as opposed to solid nanowires and polycrystalline in structure respectively Figure 5. Figure 6 represents an EDS spectrum of Co doped  $\text{TiO}_2$  nanotubes, Ti and O peaks are consistent with formation of  $\text{TiO}_2$  after thermal treatment. The Co peak was attributed to the addition of cobalt nitrate during the sol-gel process as a dopant.



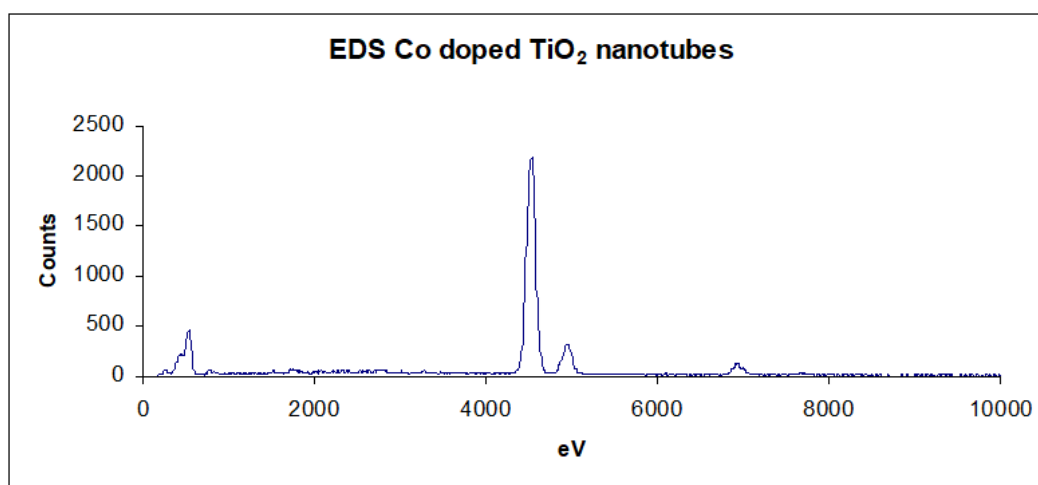
**Figure 3:** SEM image of electrospun Co- $\text{TiO}_2$ /PLA coaxial fibers.



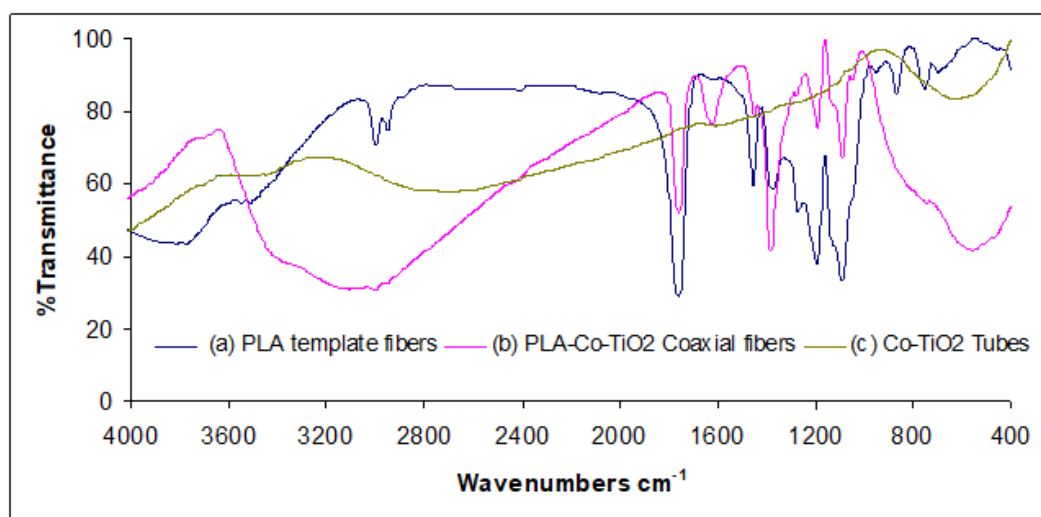
**Figure 4:** SEM image of electrospun Co doped  $\text{TiO}_2$  nanotubes.



**Figure 5:** (a) TEM image of Co doped  $\text{TiO}_2$  indicating it has a hollow interior (b) Electron diffraction pattern for Co doped  $\text{TiO}_2$  indicating polycrystallinity.



**Figure 6:** EDS spectrum for Co- $\text{TiO}_2$  nanotubes.



**Figure 7:** FTIR spectra of PLA template, PLA-Co- $\text{TiO}_2$  coaxial fibers and Co doped  $\text{TiO}_2$  tubes after thermal treatment.

Infrared spectroscopy was similarly instrumental in determination of complete elimination of the template core following thermal degradation. Figure 7(a) shows the FTIR spectrum of neat

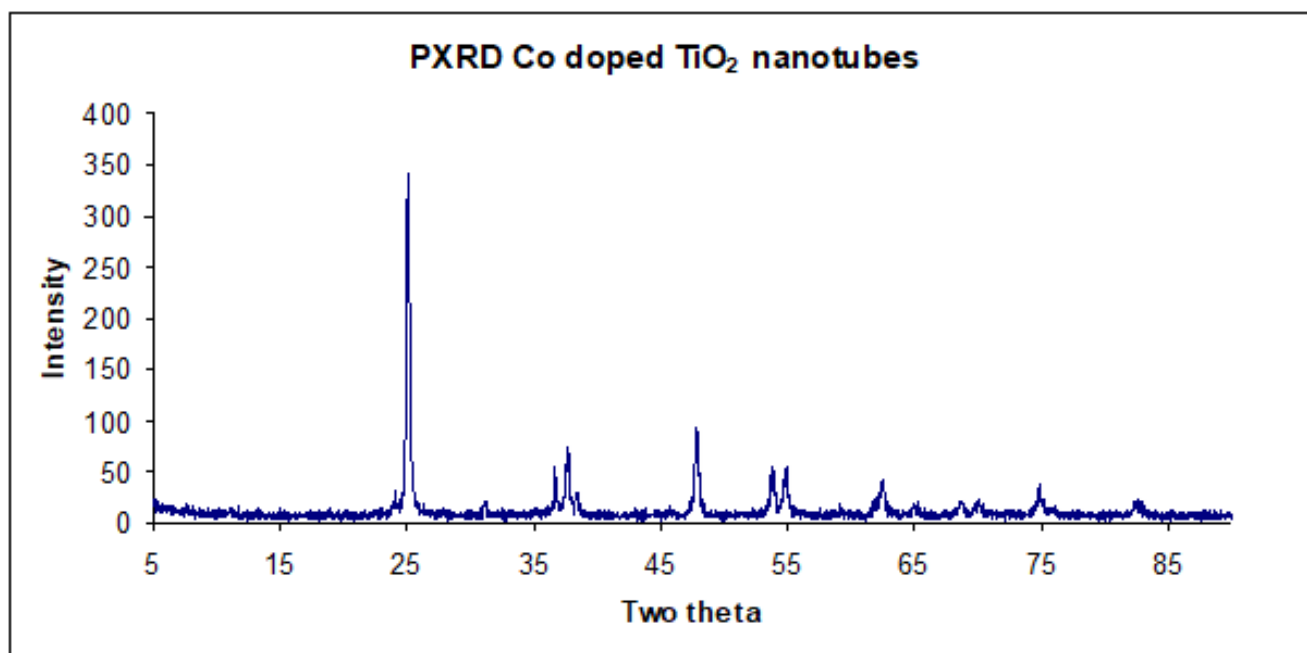
polylactide template fibers prior to sol-gel deposition. The broad absorption bands at  $\sim 3800\text{-}3500\text{cm}^{-1}$  corresponded to stretching vibrations of OH groups. The sharp bands at  $\sim 3000\text{-}2950\text{cm}^{-1}$

were assigned to the stretching vibration of aliphatic  $\text{CH}_3$  groups. The strong vibrations at  $1766\text{cm}^{-1}$  were due to  $\nu(\text{C}=\text{O})$  stretch while those at  $1458\text{cm}^{-1}$  and in the range  $1200\text{--}1000\text{cm}^{-1}$  were assigned to C-C and C-O stretching vibrations respectively [18,19]. PLA-Co-TiO<sub>2</sub> coaxial fibers FTIR spectrum, Figure 7b, exhibited similar characteristic features of the template except for the new bands at  $\sim 1624\text{cm}^{-1}$  corresponding to isopropoxide groups bonded to titanium, indicating the isopropoxide groups remain bonded during the sol-gel process. The broadened band  $\sim 500\text{--}700\text{cm}^{-1}$  was assigned to  $\nu(\text{Ti-O})$  lattice vibrations. On thermal treatment, these vibrations disappeared indicating decomposition of the polymer template material, with no sign of polylactide nor isopropoxide organic residues. The resultant broad absorption band between  $500\text{--}700\text{cm}^{-1}$  was assigned to  $\nu(\text{Ti-O})$  lattice vibrations, Figure 7(c) [20].

Structural properties and crystallite size of synthesized material was obtained by X-ray diffraction experiments. The as prepared Co-doped coaxial fibers were amorphous and on thermal treatment at  $500\text{ }^\circ\text{C}$  diffraction peaks attributed to anatase crystalline phase of TiO<sub>2</sub> were obtained [21]. The observed diffraction planes were  $2\theta=25.04^\circ(101)$ ,  $37.96^\circ(103,004,112)$ ,  $47.88^\circ(200)$ ,  $54.64^\circ(211)$ ,  $62.94^\circ(204)$ ,  $69.48^\circ(116,220)$ ,  $75.66^\circ(215)$  corresponding to the miller indices in parenthesis. A comparison with standard

diffraction spectrum, JCPDS card 21-1272 resulted in a good match with the experimental diffraction pattern. Moreover, no evidence of cobalt [22], cobalt oxide [23] or Co-Ti oxide phases, which are known to exist in bulk Co-Ti-O phase diagram were obtained [24]. This result seems to support the hypothesis that Co was distributed homogeneously in substitutional sites of anatase TiO<sub>2</sub> matrix. The average crystallite size was  $16.28\text{nm}$  as estimated using the Scherrer's equation.

Doping of TiO<sub>2</sub> with transition metal ions has been shown to influence TiO<sub>2</sub> structure depending on the concentration of the dopant. In this case, a low concentration of Co prevented it from segregating (surface nucleation) on the surface of TiO<sub>2</sub>, which inhibits sintering of amorphous TiO<sub>2</sub> particles. The absence of a mixed phase, due to effective segregation, indicates Co was homogeneously distributed within substitutional sites of the anatase matrix [15]. The influence of the dopant on the structural and textural properties of the samples can be explained based on the changes caused by the dopant on the defect structure of TiO<sub>2</sub> lattice; these changes are strongly dependent on the charge and size of the dopant ion [25,26]. Since, PXRD gives only the average crystal structure, the presence of crystal defects or impurities were determined from photoluminescence and magnetic studies (Figure 8).

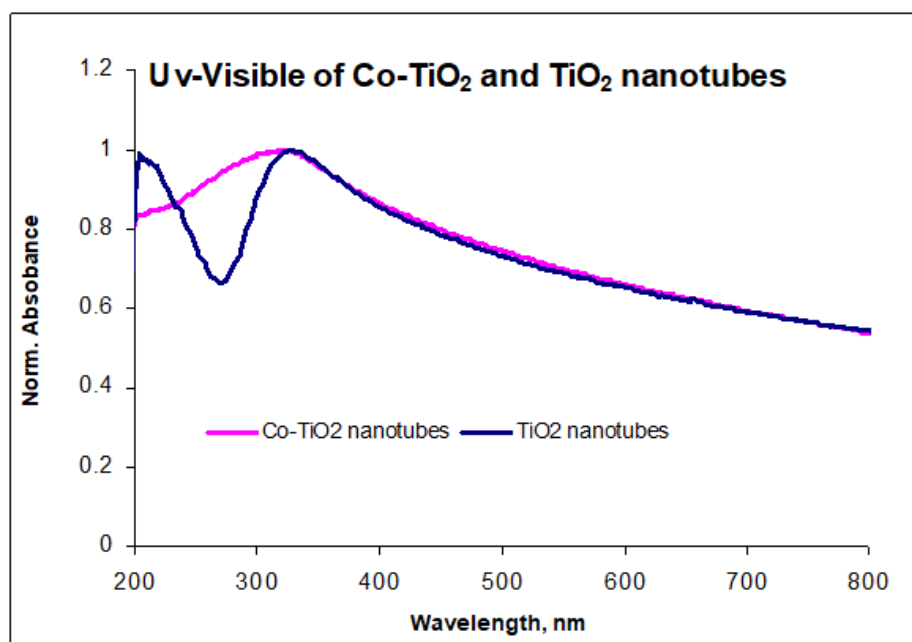


**Figure 8:** PXRD pattern of TiO<sub>2</sub> nanotubes doped with Co calcined at  $500\text{ }^\circ\text{C}$ .

### Optical and Magnetic properties

The changes in TiO<sub>2</sub> band structure were investigated using UV-visible spectroscopy. The absorption maximum for Co doped TiO<sub>2</sub> nanotubes occurred around  $325\text{nm}$ , Figure 9. This compared well with the absorption spectrum of neat TiO<sub>2</sub>, however there was no

abrupt absorption edge similar to TiO<sub>2</sub> nanotubes. The smearing of the band edge may be attributed to the presence of Co impurities within TiO<sub>2</sub> matrix [17]. Broadening results due to vibronic coupling between the valence electrons of dopant (Co) and TiO<sub>2</sub> phonons (lattice vibrations).

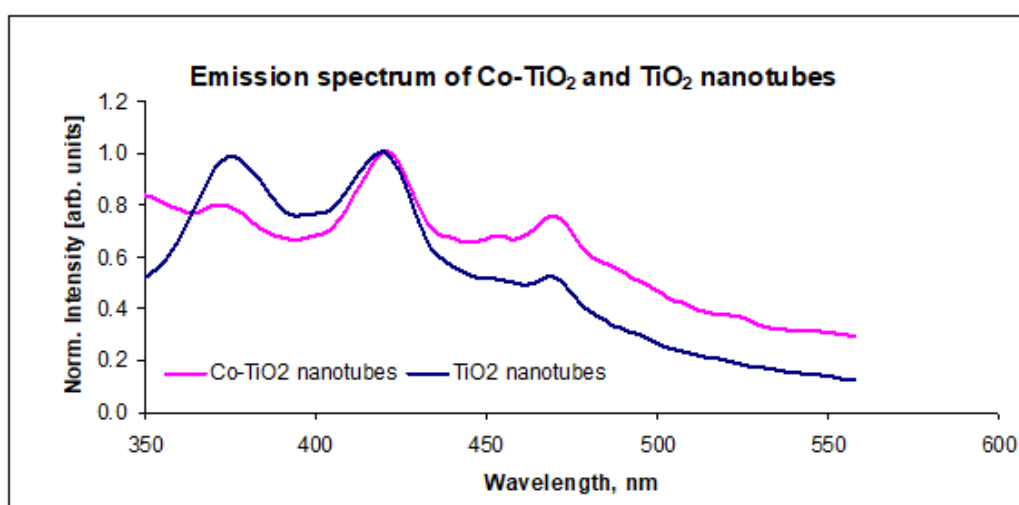


**Figure 9:** UV-visible spectrum of Co doped TiO<sub>2</sub> nanotubes.

The maximum absorption for TiO<sub>2</sub> nanotubes was similarly blue shifted relative to macro-crystalline TiO<sub>2</sub> which normally occurs around 380nm. The band gap widening was due to quantum confinement effects as a result of downsizing which is in agreement with literature studies on TiO<sub>2</sub> nanostructures [27,28]. Co doping was expected to shift the absorption edge of TiO<sub>2</sub> nanotubes to longer wavelength consistent with charge-transfer transition between the d-electrons of the dopant and the TiO<sub>2</sub> conduction band, n-type doping. The fact that the band gap did not decrease relative to TiO<sub>2</sub> nanotubes was likely due to strong band broadening, rendering precise determination of the band gap difficult, as indicated in literature for doped semiconductors [29,30]. Absorption band edge becomes more diffuse due to increasing disorder and change in lattice constant due to doping. The absence of red shift in band edge was not unique since theoretical band gap calculations for doped anatase predict the energy gap between conducting band (O 2p

and valence band (Ti 3d states) to remain unchanged on doping [31].

Photoluminescence studies provided fundamental information on energy levels lying within the band gap relative to the UV-visible absorption spectrum. The emission spectrum of Co-TiO<sub>2</sub> nanotubes was obtained by using excitation wavelength of 325nm in the range of 350-600nm at room temperature. This spectrum exhibited three emission peaks located at 374nm, 420nm and 470nm, Figure 10. The emission transition at 374nm and 420nm were attributed to highest energy direct photoemission band gap and the lowest energy indirect transition respectively [27,32]. The reduction in intensity observed in Co-TiO<sub>2</sub> nanotubes direct transition with Co doping is ascribed to the introduction of Co<sup>2+</sup> 3d states in the conduction band. This is consistent with charge-transfer transition between the d-electrons of the dopant and the TiO<sub>2</sub> conduction band, n-type doping [33].

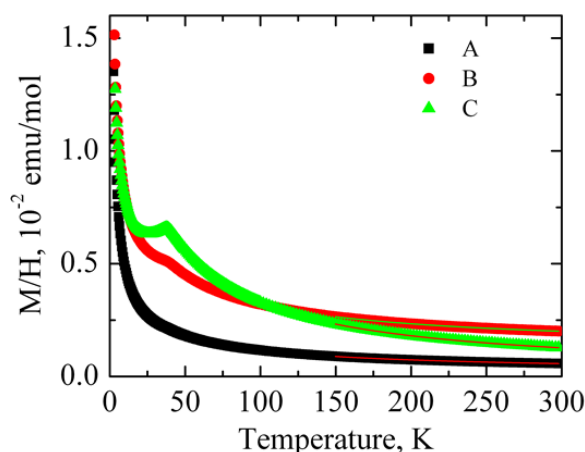


**Figure 10:** Emission spectrum of Co doped TiO<sub>2</sub> nanotubes.

The emission band at 470nm could be attributed to physical origins within the band gap: self-trapped excitons and oxygen vacancies [32,34]. For pure  $\text{TiO}_2$  nanotubes, this transition was assigned to the recombination of self-trapped excitons localized within  $\text{TiO}_2$  octahedra. The self-trapped excitons originate from band-to-band excitation where the excited electron and the remaining hole create a local deformation of  $\text{TiO}_2$  octahedra and thus localize themselves into a state in the energy gap of  $\text{TiO}_2$ . In the case of Co doped  $\text{TiO}_2$  nanotubes a combination of both factors for this emission peak (470nm), consistent with increase in peak intensity. This transition is assigned to self-trapped excitons similar to that of  $\text{TiO}_2$  nanotubes and photoluminescence due to oxygen vacancies. The oxygen vacancies are created following substitution of  $\text{Ti}^{4+}$  by  $\text{Co}^{2+}$  in the lattice in order to maintain neutrality [32,34]. Cation dopants with a charge of +4 or lower have been reported to reduce or increase the oxygen vacancy concentration in  $\text{TiO}_2$  depending on their position in the lattice. If the dopant ions ( $\text{Co}^{2+}$ ) substitute  $\text{Ti}^{4+}$  ions, oxygen vacancies tend to increase, but if they

are put in interstitial positions, the oxygen vacancy concentration decreases, lowering on crystal defects and phase transformation [15,17]. The broadening of the emission band with tailing to lower energy was attributed to broad polycrystalline size distribution and Co impurity [35].

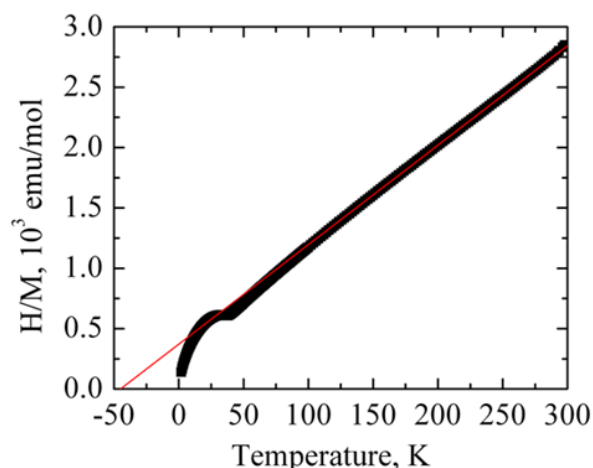
Magnetic characterization was performed on  $\text{Co-TiO}_2$  nanotubes at room temperature. The dopant concentration was similarly varied to determine the dependence of magnetization on Co content. Figure 11 shows the results of magnetic susceptibility as a function of temperature for Co doped  $\text{TiO}_2$  nanotubes. Low concentrations of Co below 4% resulted in uniform solid solution curve A, as the concentration of Co dopant increased a magnetic phase transition was observed at 38K (curve B). This was attributed to formation of  $\text{CoTiO}_3$  impurity with a Néel temperature of 38K, corresponding to antiferromagnetic-paramagnetic transition [36]. The observed transition temperature was independent of Co content and magnetic field applied consistent with formation of magnetically ordered  $\text{CoTiO}_3$  phase.



**Figure 11:** Magnetic susceptibility curve for  $\text{Co-TiO}_2$  nanotubes.

The inverse plot ( $1/\chi$  vs  $T$ ) follows Curie-Weiss behavior above 150K with  $C = 0.13$  (1)  $\text{emuK/mol}$ , Curie-Weiss temperature of  $Q = -38$  (2) K, for sample A with lowest Co content. This corresponded to 3.5 to 4.3%  $\text{Co}^{2+}$  ions with moderate antiferromagnetic exchange,

Figure 12. This compared well with the EDS estimate of  $\text{Co}^{2+}$  concentration of about 5.2-6.3% from elemental mapping of Co doped  $\text{TiO}_2$  nanotubes. For a detailed description of each sample see Table 1.



**Figure 12:** Curie-Weiss plot, indicating Curie-weiss temperature (sample A).

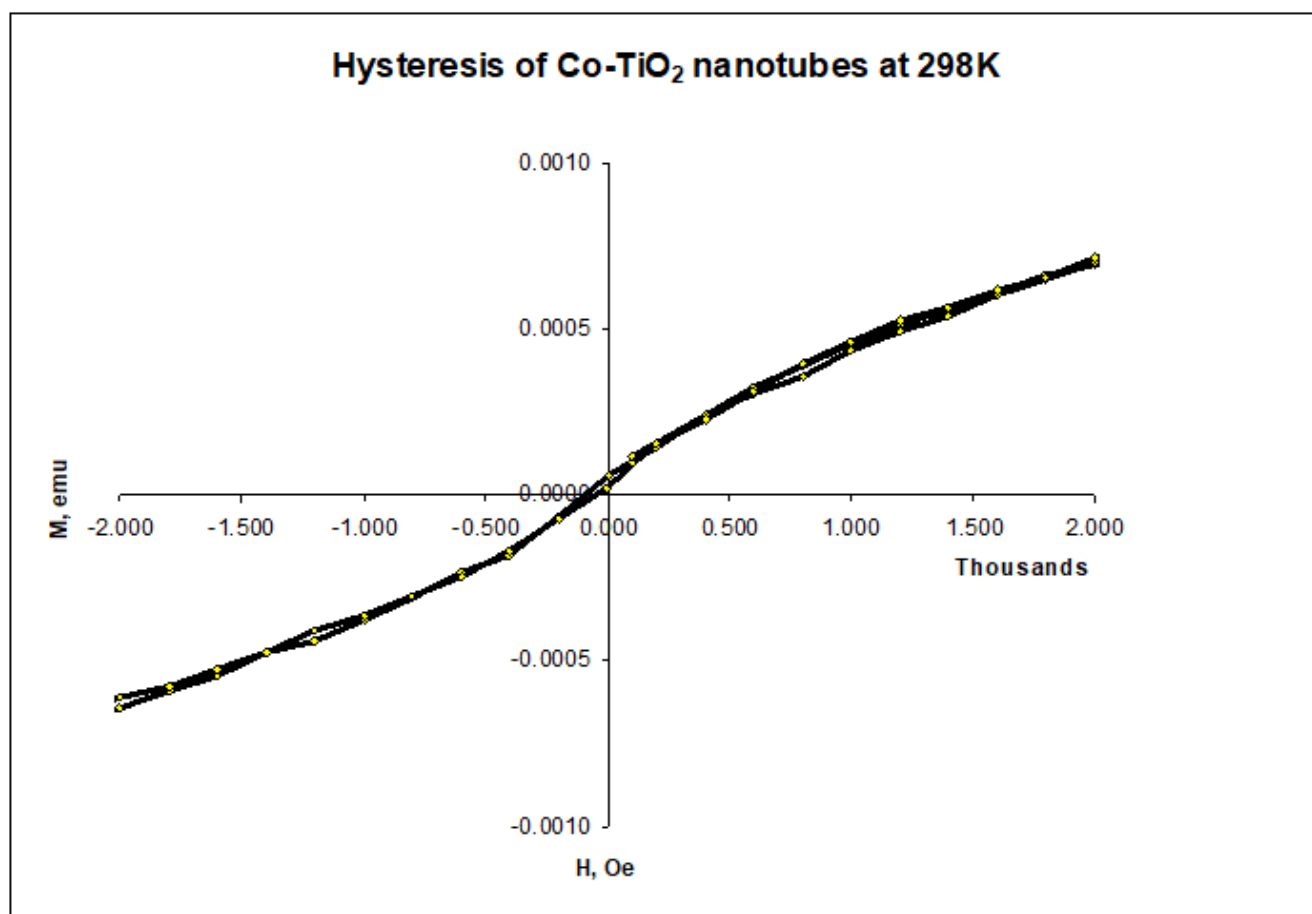


**Table 1:** Magnetic properties of Co-TiO<sub>2</sub> tubes.

Sample	38 K transition	$\chi_0$ , 10 <sup>-4</sup> emu/mol	C, emu K/mol	$\Theta$ , K	Co <sup>2+</sup> , mol.% $\mu$ eff = 3.87 $\mu$ B	Co <sup>2+</sup> , mol.% $\mu$ eff = 4.7 $\mu$ B	Co <sup>2+</sup> , mol.% $\mu$ eff = 5.2 $\mu$ B
A	No	1.8	0.13(1)	-38	6.9	4.7	3.8
B	Yes	12.6	0.24(1)	-24	12.8	8.7	7.1
C	Yes	1.2	0.36(1)	-15	19.2	13	10.7

Figure 12 shows 298K hysteresis loop measured for Co-TiO<sub>2</sub> (A), without magnetic phase transition. The linear dependence of magnetization on magnetic field and the absence of a magnetic hysteresis loop indicates Co doped TiO<sub>2</sub> nanotubes prepared were paramagnetic at lower percent loadings. The absence of

room temperature ferromagnetism in Co doped TiO<sub>2</sub> nanotubes may be attributed to calcination in air which may have favoured formation of larger fraction of antiferromagnetically coupled Co<sup>2+</sup> ions resulting in negligible contribution to overall magnetization [36,37] Figure 13.

**Figure 13:** Hysteresis loop for Co-TiO<sub>2</sub> nanotubes, sample A.

## Conclusion

Co doped TiO<sub>2</sub> submicron tubes were fabricated using electro-spun polymer fibers as templates via Sol-gel deposition. Thermal degradation of the template core at 500 °C resulted in polycrystalline anatase TiO<sub>2</sub> tubes with Co distributed within substitutional sites. UV-visible measurements indicated band gap widening due to quantum confinement effects, while photoluminescence experiments confirmed presence of point defects due to oxygen vacancies. Magnetization measurements indicated Co doped TiO<sub>2</sub> nanotubes were paramagnetic with a magnetic phase transition occurring reported in samples with higher than 4% Co loading.

Future studies will focus on improving magnetic ordering and application these structures in spintronics.

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